

massless particle: only possible to boost to a frame with

$$p^\mu = (E, \underbrace{0, 0, 0}_\text{little group is SO(2)}, E) \quad \text{because } p^\mu p_\mu = -\frac{E^2}{c^2} + |\mathbf{p}|^2 = -m^2 c^2 = 0$$

only one generator, J_{12} , eigenvalue is called helicity λ . Like spin s , λ is a half integer.

complete state $|p_\mu, \lambda\rangle$

examples: massive:

Higgs mechanism	$S=0$	Higgs boson
	$S=1/2$	electron
	$S=1$	W and Z bosons

massless:

$\lambda=0$	KG field for $m=0$
$\lambda=1$	photon $\Rightarrow$ <u>electromagnetic field</u>
$\lambda=2$	graviton $\Rightarrow$ gravity

Question: What are the Lagrangians for all these theories?

5. Electrodynamics

is governed by Maxwell equations for electric field $\vec{E}$ and magnetic field $\vec{B}$

$$1) \nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \quad \text{how to compute?}$$

$$\vec{\nabla} = \vec{e}_1 \partial_1 + \vec{e}_2 \partial_2 + \vec{e}_3 \partial_3, \text{ and same for } \vec{E} = \vec{e}_1 E_1 + \vec{e}_2 E_2 + \vec{e}_3 E_3, \vec{B} = \vec{e}_1 B_1 + \vec{e}_2 B_2 + \vec{e}_3 B_3$$

$$\vec{e}_i \times \vec{e}_j = \sum_k \epsilon_{ijk} \vec{e}_k \quad \leadsto \text{for example}$$

$$\partial_1 E_2 - \partial_2 E_1 = -\frac{1}{c} \frac{\partial B_3}{\partial t} \quad \dots$$

$$2) \nabla \cdot \vec{B} = 0, \quad 3) \nabla \cdot \vec{E} = \rho$$

$$\vec{e}_i \cdot \vec{e}_j = \delta_{ij}$$

$$\partial_1 B_1 + \partial_2 B_2 + \partial_3 B_3 = 0$$

$$4) \nabla \times \vec{B} = \frac{1}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t} \quad (\text{all in Heaviside-Lorentz unit system})$$

1) + 2) homogeneous, solved for $\vec{E} = \vec{B} = 0$

3) + 4) inhomogeneous due to charge ρ and current $\vec{j}$

5.1. Vector and electric potential

$$2) \text{ is solved by } \vec{B} = \vec{\nabla} \times \vec{A}$$

$$\text{because } \vec{\nabla} \cdot (\vec{\nabla} \times \vec{A}) = 0 \quad \text{grad}(\text{curl}(\vec{A})) = 0$$

$$1) \text{ then implies } \vec{\nabla} \times \left(\vec{E} + \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right) = 0$$

$$\Rightarrow \vec{E} = -\frac{1}{c} \frac{\partial \vec{A}}{\partial t} - \nabla \phi \quad \sim \text{is allowed because}$$

$$\nabla \times (\nabla \phi) = 0$$

$$\text{curl}(\text{div } \phi) = 0$$

5.2. Covariant reformulation

$$A^\mu = (\phi, A_1, A_2, A_3), \quad A_\mu = (-\phi, A_1, A_2, A_3)$$

for which we compute $X^\mu = (ct, x^1, x^2, x^3)$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu = 2 \partial_{[\mu} A_{\nu]}$$

independent components: F_{0i} and F_{ij} ($i, j = 1, 2, 3$)

$$F_{0i} = \frac{\partial A_i}{\partial x^0} - \frac{\partial A_0}{\partial x^i} = \frac{1}{c} \frac{\partial A_i}{\partial t} + \frac{\partial \phi}{\partial x^i} = -E_i //$$

$$F_{ij} = \partial_i A_j - \partial_j A_i = \sum_k \epsilon_{ijk} B_k$$

$$\Rightarrow F_{12} = B_3, \quad F_{23} = B_1, \quad F_{13} = -B_2$$

Putting all together we get

$$F_{\mu\nu} = \begin{pmatrix} 0 & -E_1 & -E_2 & -E_3 \\ E_1 & 0 & B_3 & -B_2 \\ E_2 & -B_3 & 0 & B_1 \\ E_3 & B_2 & -B_1 & 0 \end{pmatrix} \quad \underline{\text{field strength tensor}}$$

Question: How do the Maxwell equations look in this formulation?

1)+2) $\nabla \cdot (\nabla \times \vec{A}) = 0$ and $\nabla \times (\nabla \cdot \phi) = 0$
are equivalent to

$$\underline{\underline{\partial_\lambda (\partial_\mu A_\nu - \partial_\nu A_\mu) + \partial_\mu (\partial_\nu A_\lambda - \partial_\lambda A_\nu) + \partial_\nu (\partial_\lambda A_\mu - \partial_\mu A_\lambda) = 0}}$$

or equivalently $\boxed{3 \partial_{[\mu} F_{\nu\sigma]} = 0}$ Bianchi identity

3+4) $F^{\mu\nu} = \eta^{\mu\alpha} \eta^{\nu\beta} F_{\alpha\beta}$

with $F^{\mu\nu} = -F^{\nu\mu}$, $F^{0i} = -F_{0i}$
 $F^{ij} = F_{ij}$

$\boxed{\partial_\nu F^{\mu\nu} = \frac{1}{c} j^\mu}$ field equation with

$$j^0 = c \partial_\nu F^{0\nu} = c (\partial_1 E_1 + \partial_2 E_2 + \partial_3 E_3) \\ = c \nabla \cdot \vec{E} = c \rho$$

$$\partial_\nu F^{i\nu} = -\frac{1}{c} \frac{\partial}{\partial t} E_i + \sum_{j,k} \partial_j B_k \epsilon_{jki} \\ = \left(-\frac{1}{c} \frac{\partial \vec{E}}{\partial t} + \nabla \times \vec{B} \right)_i = \frac{1}{c} j_i$$

$\Rightarrow \boxed{j^\mu = (c\rho, j_1, j_2, j_3)}$ current four vector

in vacuum $j^\mu = 0$

Bonus: We can now easily define electrodynamics in dimension other than 4.

5.3. Action

$$S = \frac{1}{4} \int d^4x \, F_{\mu\nu} F^{\mu\nu}$$

vacuum action

check:

$$\begin{aligned} \delta S &= \frac{1}{2} \int d^4x \, \delta F_{\mu\nu} F^{\mu\nu} = \int d^4x \, \partial_\mu \delta A_\nu F^{\mu\nu} \\ &= \int d^4x \, \delta A_\nu \underbrace{\partial_\mu F^{\mu\nu}}_{=0} = 0 \end{aligned}$$

e.o.m : $\partial_\nu F^{\mu\nu} = 0$